

Simulating CME initiation and evolution: state-of-the-art

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Space weather : solar drivers

- drivers are of **solar origin**: viz. transient phenomena superposed on the solar wind:
 - CMEs (most prominent)
 - eruptive flares
 - SEPs
 - ...
- basic physical mechanisms **not fully understood**
- predictions and forecasts '**can be improved**'

CMEs :

- typ. 400 km/s , $10^{12} - 10^{13} \text{ kg!}$
- $E = 10^{24} - 10^{25} \text{ Joule}$
- known since 30 yrs only!
- **they play a crucial role in SW!**

Solar wind modeling : recent developments

- CSEM (Univ. Michigan): **first coupled model** of the inner heliosphere extending from the low solar corona to the well beyond Earth orbit
- ⇒ Space Weather Modeling Framework (SWMF): a high-performance flexible computational tool that *enables coupling state-of-the art models of the solar corona, the solar wind and solar energetic particles*
(Gombosi et al. (2004), Tóth et al. (2005))

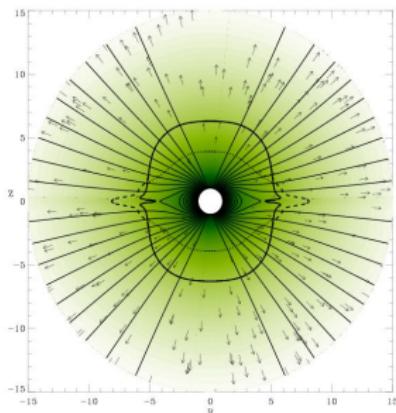
Other groups also developed advanced solar wind models, e.g.:

- Lionello et al. (2003): **3D MHD model of the solar corona and of the solar wind** (polytropic eq. improved by incorporating thermal conduction ($\parallel \mathbf{B}$), radiation losses, and heating into the energy eq.)

Recent developments

- Lee et al. (2004): extended the CISM heliospheric model CORHEL to 10 AU to investigate how well this **solar magnetogram-based** 3D MHD model describes the solar wind influence on Saturn's magnetosphere
- Odstrcil et al. (2004): coupled numerical wind model where
 - (1) the ambient solar wind is derived from coronal models utilizing **photospheric magnetic field observations** and
 - (2) transient disturbances are derived from geometrical and kinematic **fitting of coronagraph observations** of coronal mass ejections (CMEs)

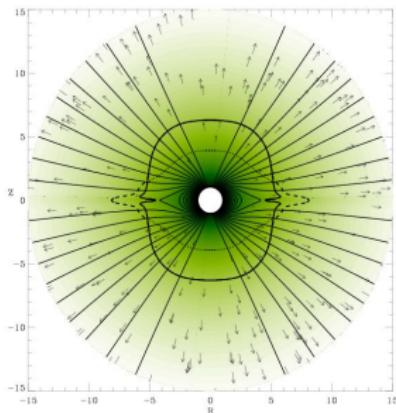
Solar wind models @ CPA



Polytropic Wind

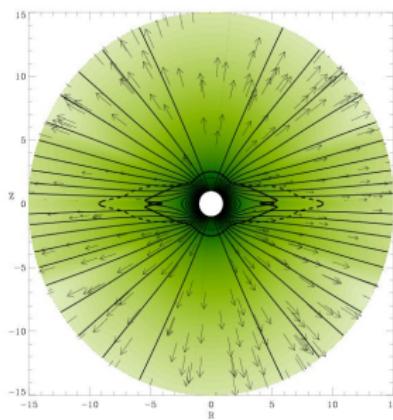
Color: density (log-scale), black lines:
magnetic field lines, arrows: velocity

Solar wind models @ CPA



Polytropic Wind

Color: density (log-scale), black lines: magnetic field lines, arrows: velocity

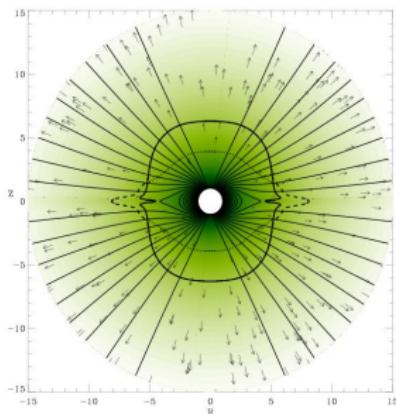


'Manchester' wind

Extra heating source term:

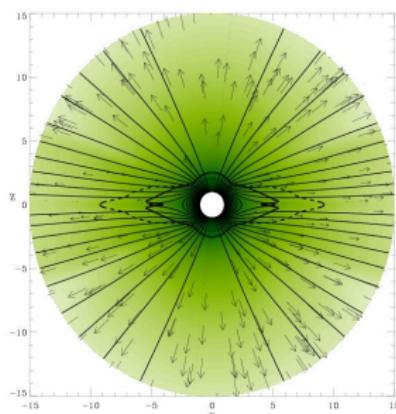
$$Q = \rho q_0 e^{-\frac{(r-r_0)^2}{\sigma^2}} \left(T_0 - \gamma \frac{p}{\rho} \right)$$

Solar wind models @ CPA



Polytropic Wind

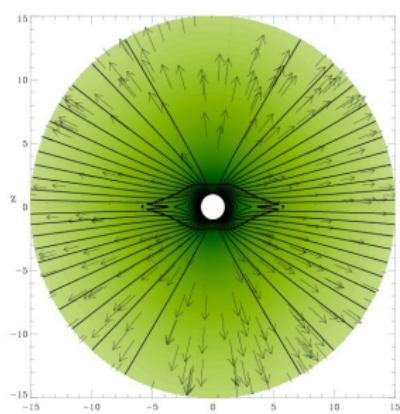
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Polytropic Wind with
Alfvén Waves

Has additional Alfvén wave pressure gradient.

Wind characteristics at $30R_{\odot}$

	Model 1	Model 2	Model 3
<i>Density</i> [m ⁻³]			
Pole	5.6×10^8	1.02×10^9	3.08×10^9
Equator	7.27×10^8	1.83×10^9	2.87×10^9
Ratio	0.77	0.56	1.07
<i>Velocity</i> [km/s]			
Pole	323	727	675
Equator	293	358	374
Ratio	1.1	2.03	1.8
<i>Temperature</i> [K]			
Pole	0.82×10^6	1.13×10^6	0.89×10^6
Equator	0.83×10^6	0.29×10^6	0.89×10^6
Ratio	0.99	3.87	1.0
<i>Magnetic field</i> [G]			
Pole	6.04×10^{-4}	3.7×10^{-4}	3.9×10^{-4}
Equator	6.1×10^{-5}	1.2×10^{-4}	2.0×10^{-4}
Ratio	9.89	3.06	1.95

Recent developments

- heliospheric models of CME propagation and evolution provide an **important insight into the dynamics of CMEs** and are a **valuable tool for interpreting interplanetary in situ observations**
- Linker et al. (2003): **need to study coronal initiation and solar wind propagation together**
⇒ serious efforts have been undertaken to reach this goal!
- Roussev et al. (2003a,b) developed the capability to use observed **synoptic magnetograms** to drive the coupled corona-solar-wind model and to simulate the SW. The tool was used to simulate a 3D flux-rope model for a CME **based on a loss of equilibrium**, i.e. not on an initially unsatisfied force balance as in many earlier (and current) simulations!

Recent developments (cont.)

- Manchester et al. (2004a,b) also modeled erupting flux ropes and the resulting CMEs in **full 3D magnetohydrodynamics**
- Sokolov et al. (2004) then included a field line advection model and this coupled corona-solar-wind-SEP model thus **includes the energetic particle environment** at any point in the inner heliosphere in the simulations
- Manchester et al. (2005) focussed on the **CME shock and shear structures** relevant for particle acceleration, while Lugaz et al. (2005) concentrated on the **evolution of the density structure of the CMEs**

Recent developments (cont.)

- Jacobs et al. (2005) also focussed on the propagation of fast-CME generated MHD shock waves and made a first attempt to **quantify the effect of the background solar wind model** on this evolution by superposing the same simple CME model on three different 2.5D (axisymmetric) wind models
- Chané et al. (2005) then studied the **effect of the CME initiation parameters on the CME evolution**, in particular the polarity of the initial magnetic flux rope \Rightarrow *affects velocity, path, spread angle, etc.*

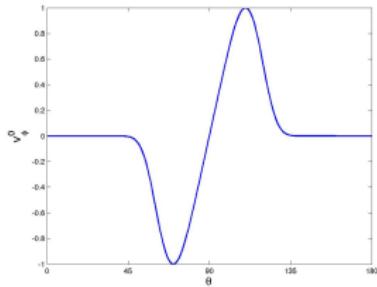
CME initiation: foot point shearing

Add extra azimuthal velocity

v_ϕ^0 at the solar surface to
shear the footpoints of the
magnetic field lines

$$v_\phi^0 = v_0(t)\Theta e^{(1-\Theta^4)/4}$$

$$\text{with } \Theta = \frac{\theta - \pi/2}{\Delta\theta_m}$$



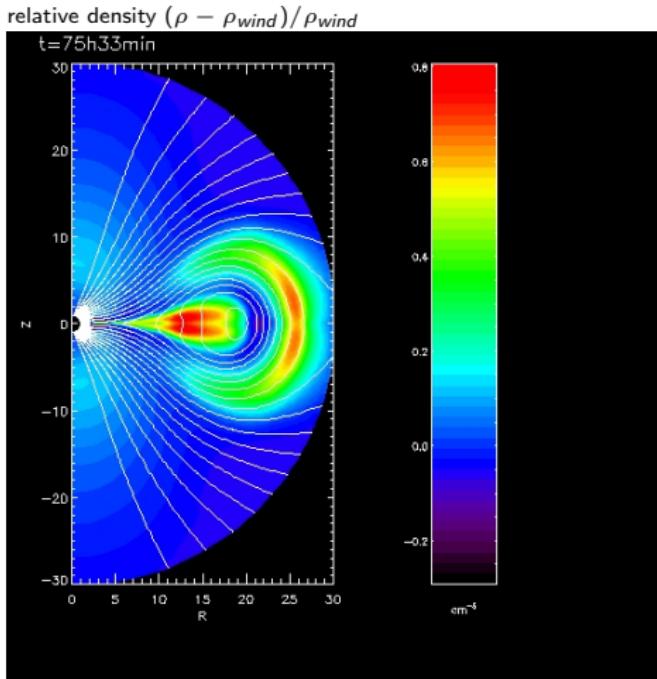
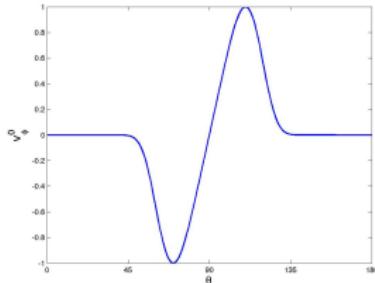
Background wind: **Model 1**, maximum shear velocity: **6 km/s**.

CME initiation: Shearing

Add extra azimuthal velocity v_ϕ^0 at the solar surface to shear the footpoints of the magnetic field.

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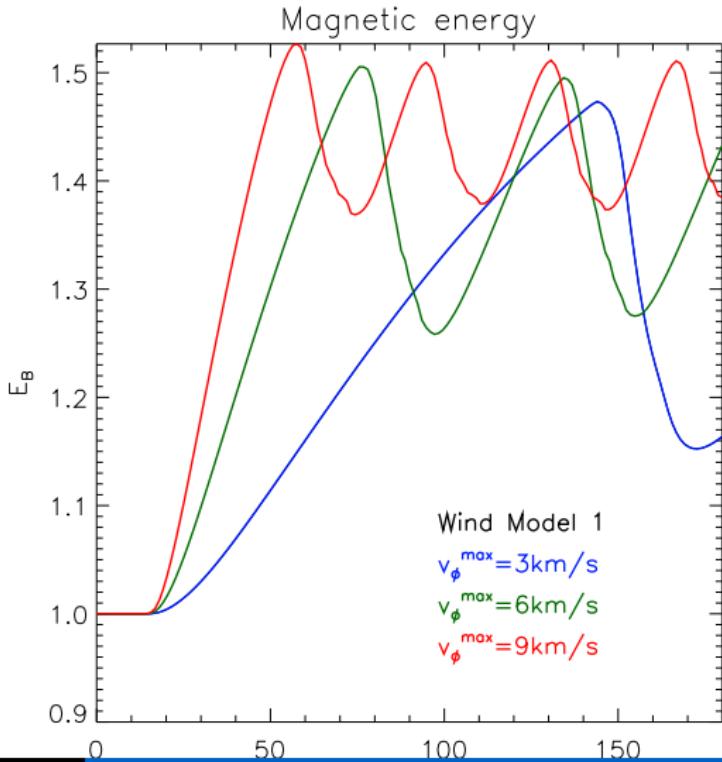


Background wind: **Model 1**, maximum shear velocity: **6 km/s**.

CME initiation: Shearing - parameter studies

Shearing rate & background wind both affect

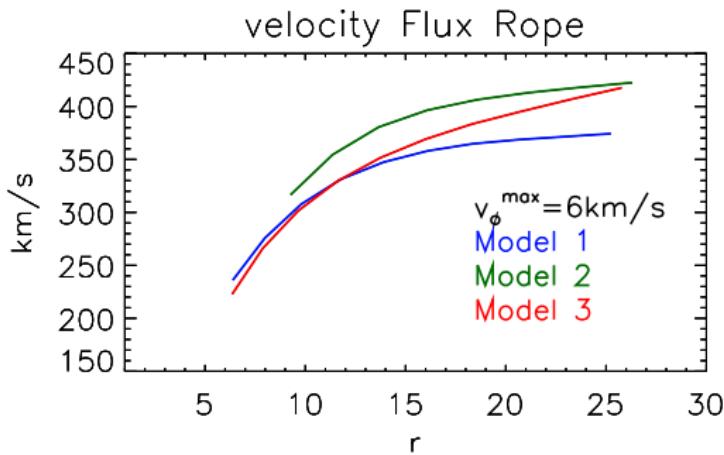
- Δt to reach instability
- instability threshold in terms of energy
- amount of energy released
- velocity/acceleration of flux rope
- ...



CME initiation: Shearing - parameter studies

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Magnetic helicity

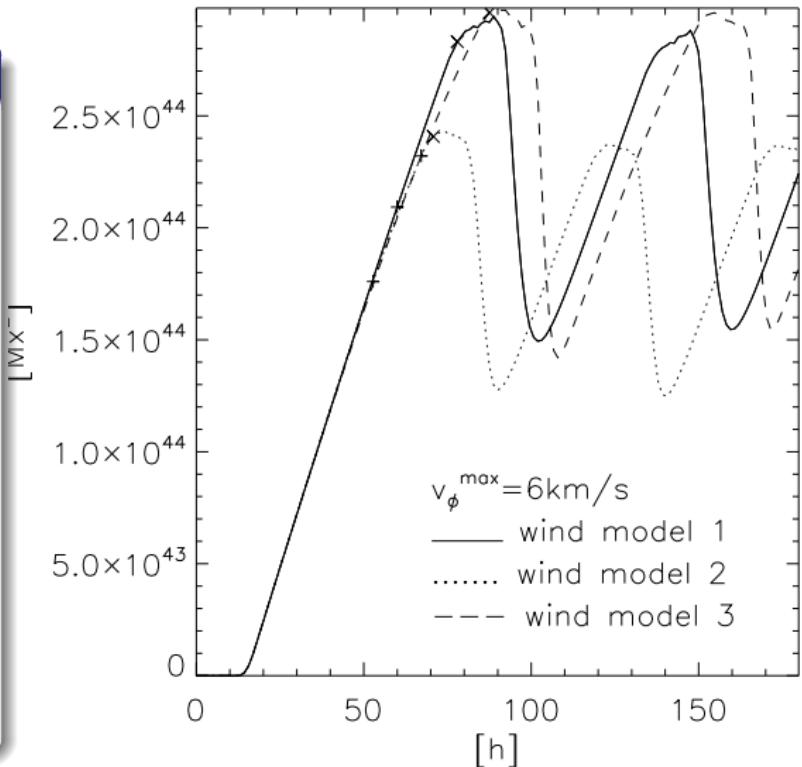
$$H_r = \int_V (\mathbf{A} + \mathbf{A}_p) \cdot (\mathbf{B} - \mathbf{B}_p) dV$$

simplifies to

$$H_r = 2 \int_V A_\varphi B_\varphi dV$$

in axi-symmetric case

- ⇒ **easy to compute**
- ⇒ **difficult to measure!**
- ⇒ **helicity injection rate**



Magnetic helicity

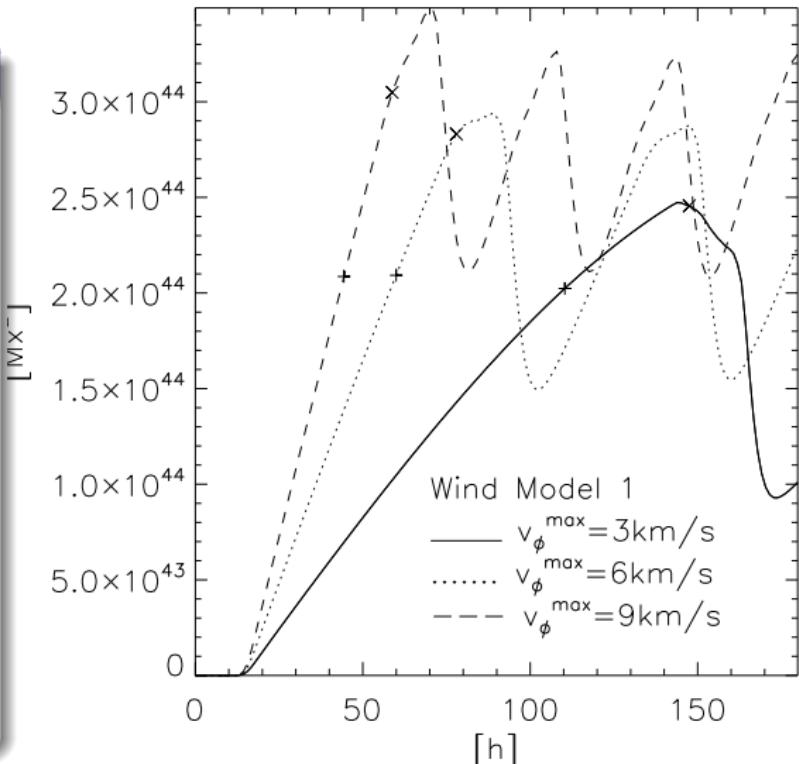
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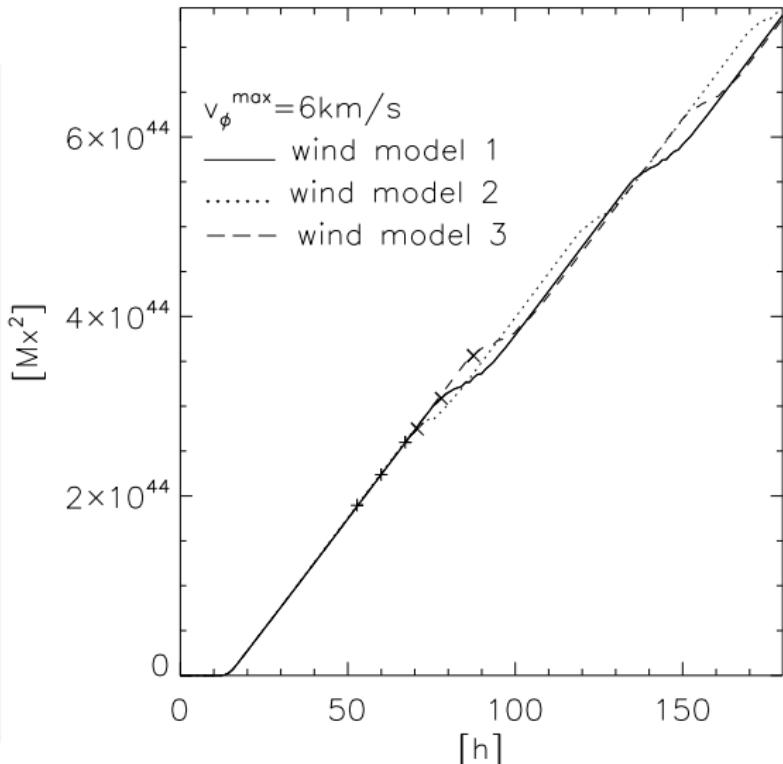
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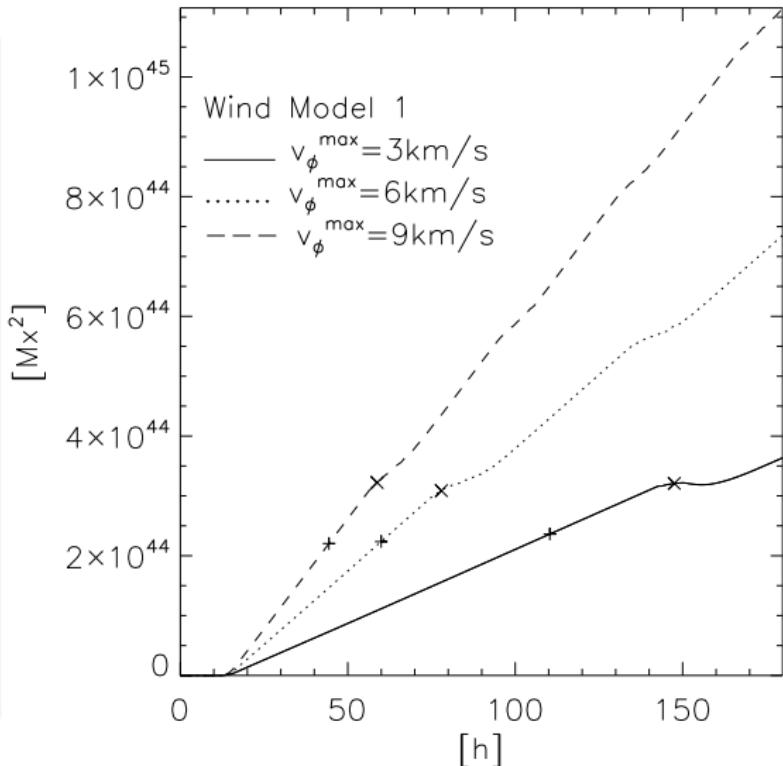
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CME initiation: flux emergence / cancellation

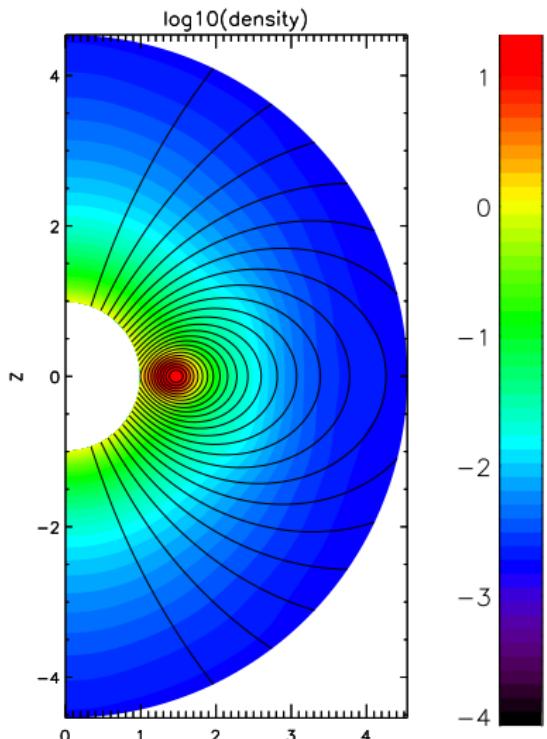
Initial model

- cf. Chen & Shibata (2000)
- + **physics** : MHD (incl. gravity)
- + **geometry** : 2D (axisymmetric)
- + dipole field

OR

+ **solar wind** \Rightarrow

$$\vec{B}_0 = \vec{B}_{LC} + \vec{B}_{IC} + \vec{B}_{BG}$$

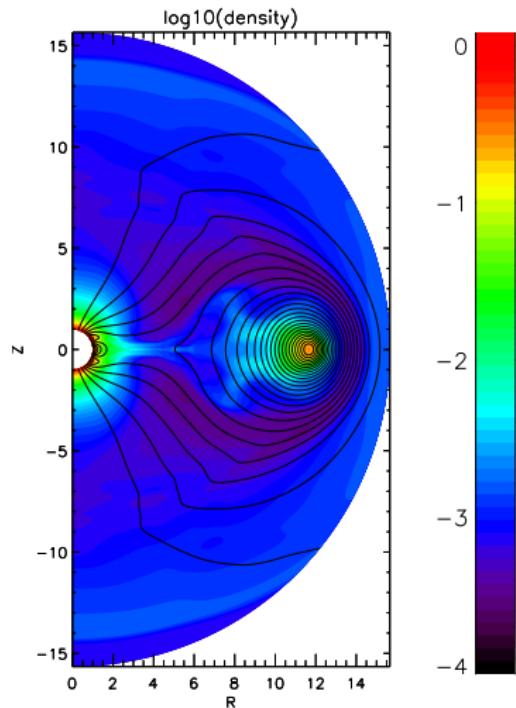
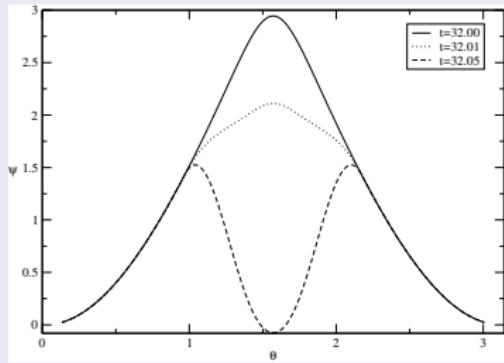


CME initiation: flux emergence / cancellation

Addition of flux

cf. Forbes & Priest ('84), Chen & Shibata ('00)

- at lower boundary ($r = 1R_{\odot}$)
- in region $\frac{\pi}{2} - 0.6 \leq \theta \leq \frac{\pi}{2} + 0.6$
- BC: $A_{\varphi} = A_{\varphi}(t_0) + c_e A_{\varphi}^+ \frac{t - t_0}{t_e - t_0}$



Case A : varying flux emergence rate, fixed c_e (total flux)

Parameter study : t_e

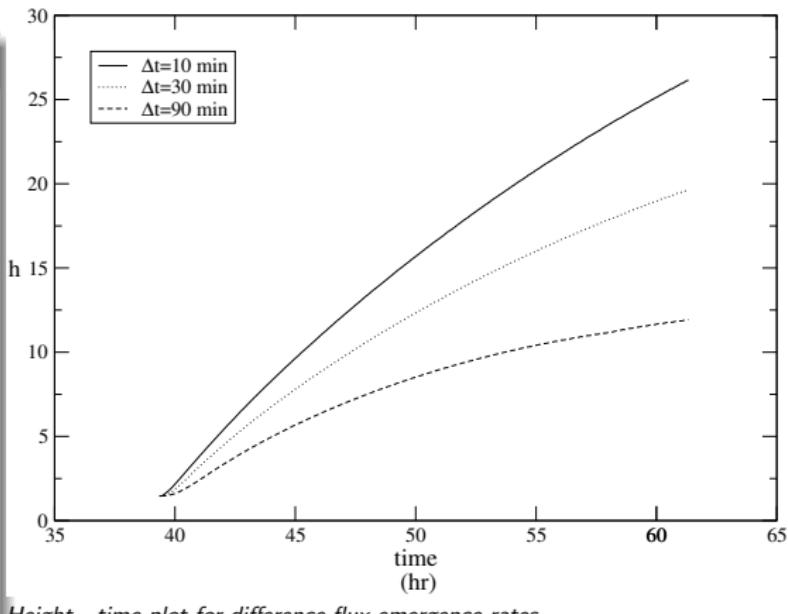
- BC:

$$A_\varphi = A_\varphi(t_0) + c_e A_\varphi^+ \frac{t - t_0}{t_e - t_0}$$

- fixed amount of flux:

$$2\pi c_e \psi_0 \approx -6.6 \times 10^{22} \text{ Mx in Northern hemisphere}$$

- vary *flux emergence rate*, i.e. $2\pi c_e \psi_0 / \Delta t$ from -1.10×10^{20} to $-1.22 \times 10^{19} \text{ Mx/s}$



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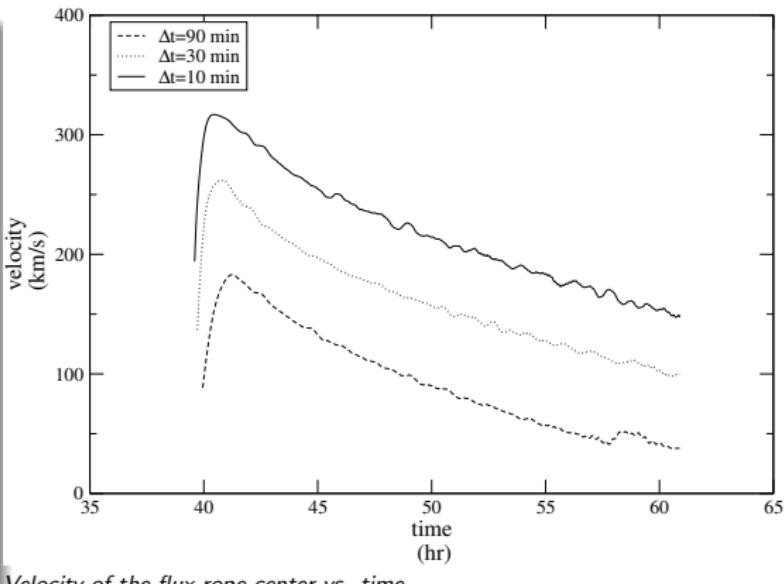
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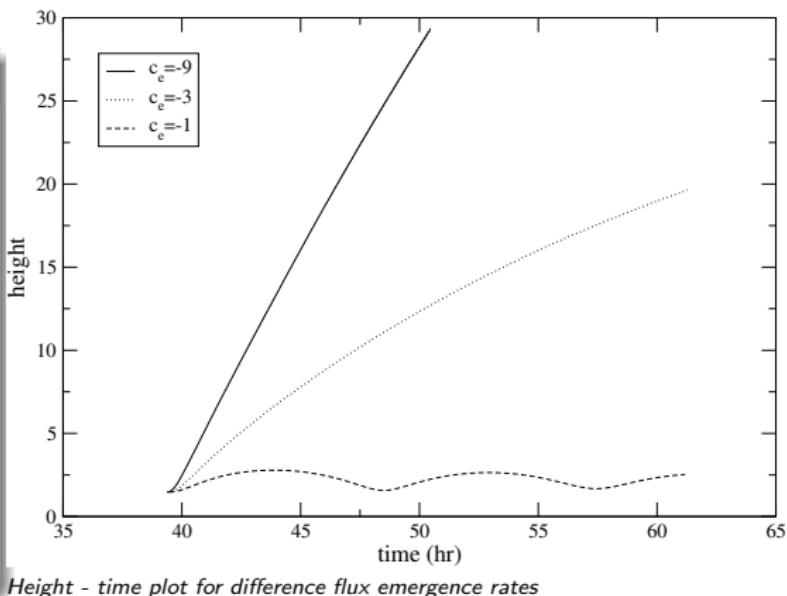
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Case B : varying flux emergence rate, fixed t_e (Δt)

Parameter study : t_e

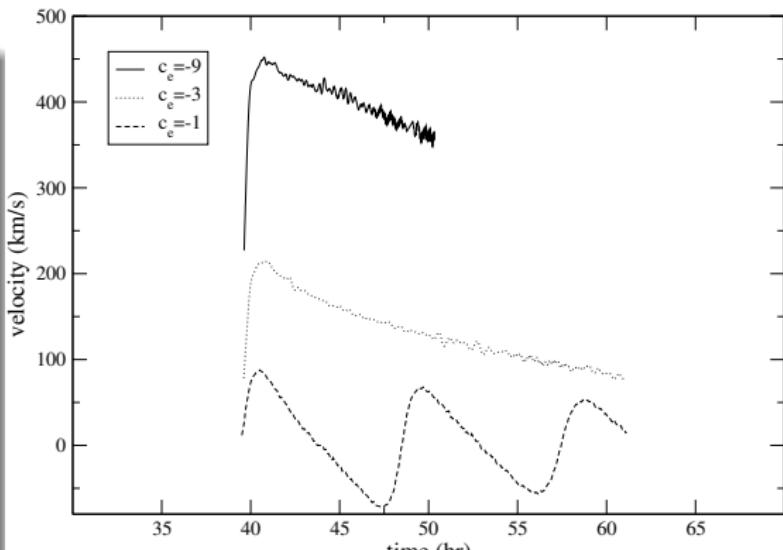
- fixed $\Delta t = 30$ min
- total flux* varies from -1.1×10^{22} Mx ($c_e = -1$) to -9.9×10^{22} Mx ($c_e = -9$)
- flux emergence rate* varies from -1.22×10^{19} to -1.10×10^{20} Mx/s



Case B : varying flux emergence rate, fixed t_e (Δt)

Parameter study : t_e

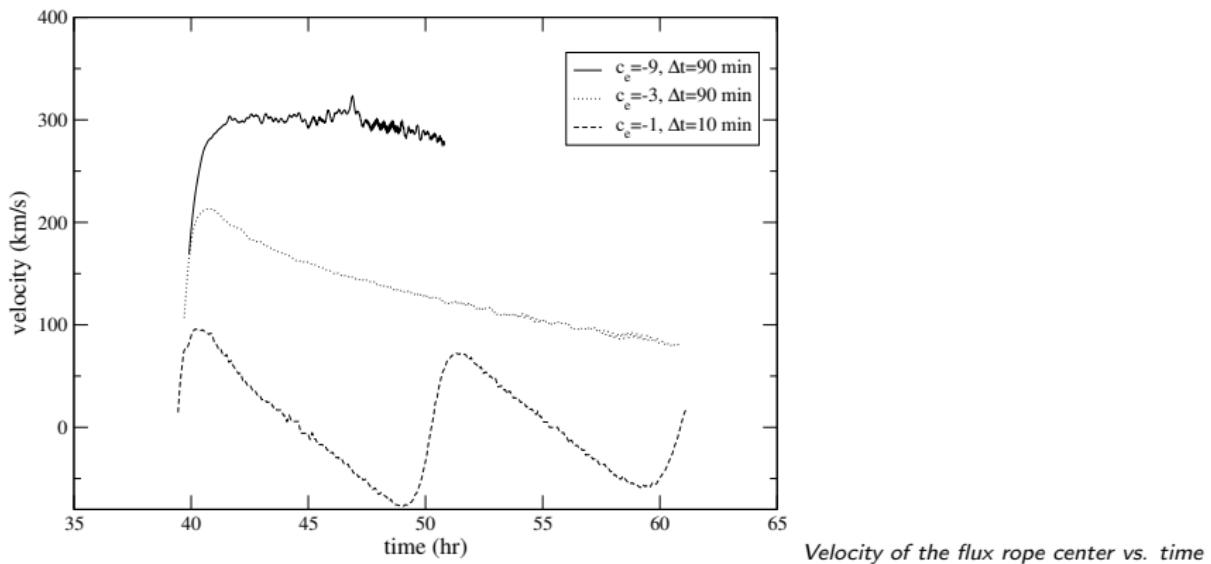
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- flux emergence rate* varies from -1.22×10^{19} to -1.10×10^{20} Mx/s



Velocity of the flux rope center vs. time

Case C : fixed flux emergence rate, varying t_e and c_e

	$\Delta t =$	10 min	30 min	90 min
$c_e =$	-1	3.67×10^{19} Mx/s	1.10×10^{20} Mx/s	
	-3	1.22×10^{19} Mx/s	3.67×10^{19} Mx/s	1.10×10^{20} Mx/s
	-9		1.22×10^{19} Mx/s	3.67×10^{19} Mx/s



CME evolution : recent developments

- Aran et al. (2004, 2005) developed and applied a tool for rapid **predictions of proton flux and fluence profiles** observed during gradual solar energetic particle (SEP) events, for the upstream part of the shock (SOLPENCO)
 - ⇒ contains data base with set of IP scenarios for SEPs, defined by the
 - (1) *heliolongitude* of the parent solar activity (from E75 to W90)
 - (2) *position of the observer* (either at 0.4 AU or 1.0 AU) and
 - (3) *initial shock velocity* (between 750 km/s and 1800 km/s at 18 solar radii) of the simulated shock

Recent developments (cont.)

- Fry et al. ('03,'04,'05) used *1st generation, real-time, combined "quiet" and "event-driven" solar wind model* to examine more than 600 solar flares since 1997 (rise, maximum, and decline of Solar Cycle 23) that were accompanied by metric radio Type II radio bursts and/or by partial or full halo CMEs...
⇒ experience in forecasting the flare-shocks' times of arrival (and, by implication, the following ICMEs), was summarized by discussing several specific cases, including the April/May 2001 and Halloween 2003 extreme events (Sun et al. ('03), Dryer et al. ('04), Wu et al. ('05a))

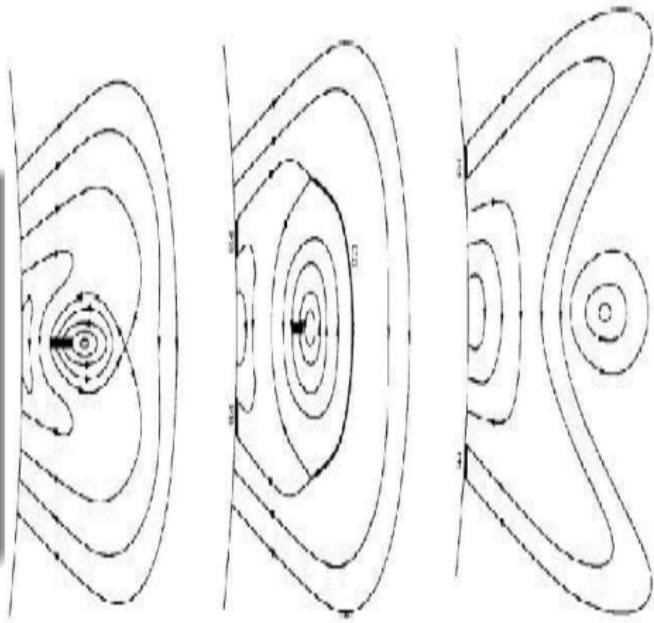
Recent developments (cont.)

- Wu et al. ('05b) also developed a **data driven 3D MHD model** to investigate the non-potentiality of active-region corona (innovative t -dependent projected characteristic BCs)
 - Wu et al. ('04) examined the Low and Zhang (2002) suggestion that the two types of CMEs (i.e. constant speed (fast) and accelerated (slow)) are caused by the initial magnetic topology
- ⇒ simulation shows that, in addition to the magnetic topology, the *solar surface condition also plays an important role to determine the two types of CMEs*

CME evolution up to $30 R_{\odot}$: creating shocks

Simple CME model

- evolution normal/inverse CMEs
(Low & Zhang (2003))
confirmed
- effect SW quantified
- initial magnetic polarity affects
evolution path

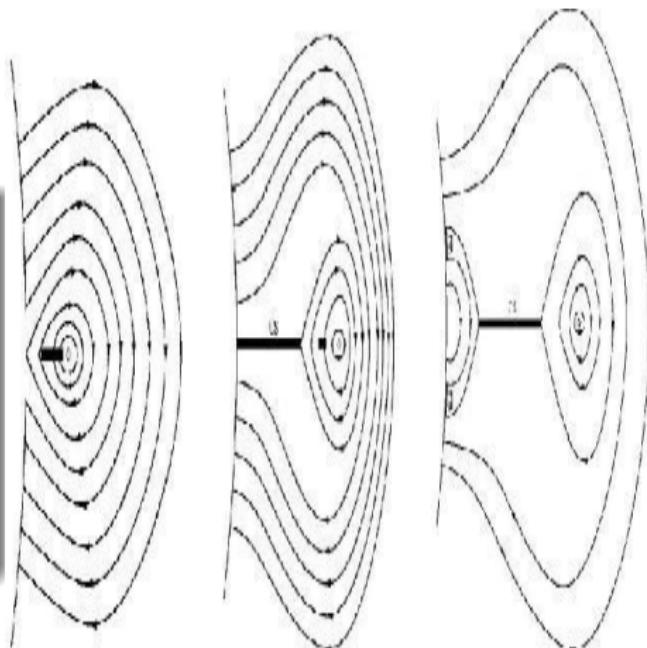


Evolution of a 'normal' CME

CME evolution up to $30 R_{\odot}$: creating shocks

Simple CME model

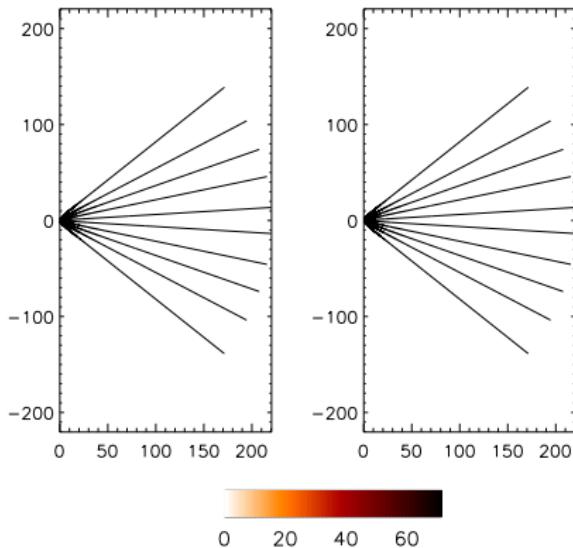
- evolution normal/inverse CMEs
(Low & Zhang (2003))
confirmed
- effect SW quantified
- initial magnetic polarity affects
evolution path



Evolution of an 'inverse' CME

CME evolution up to 1 AU

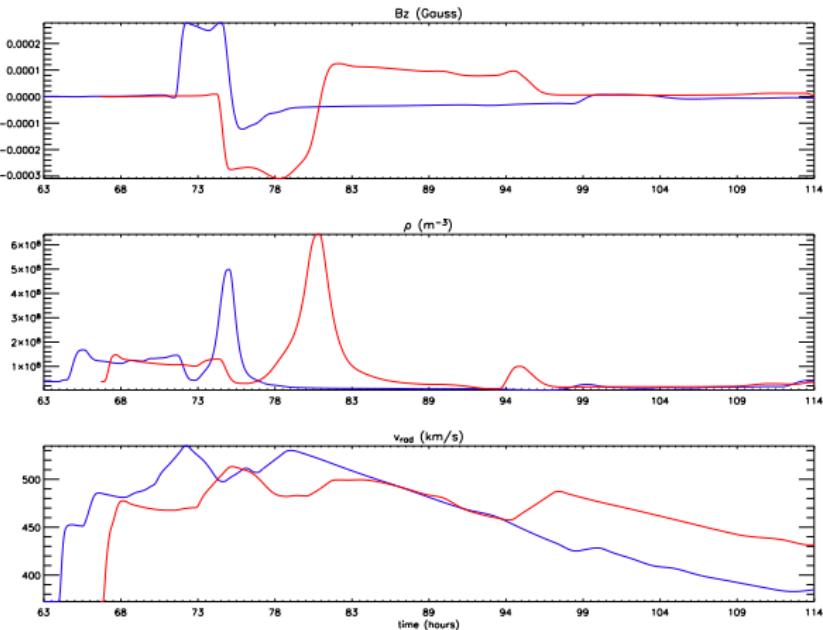
- self-similar evolution stops beyond $30 R_{\odot}$
- difference normal/inverse polarity much smaller
(e.g. density distr.)
- higher wind density at equator leads to serious deformation (compression) of the CMEs
- only difference :
 - about 6 hrs \neq in arrival time
 - orientation of field



Relative density and magnetic field lines (INVERSE - NORMAL)

Simulated satellite data at 1 AU (Wind model 2)

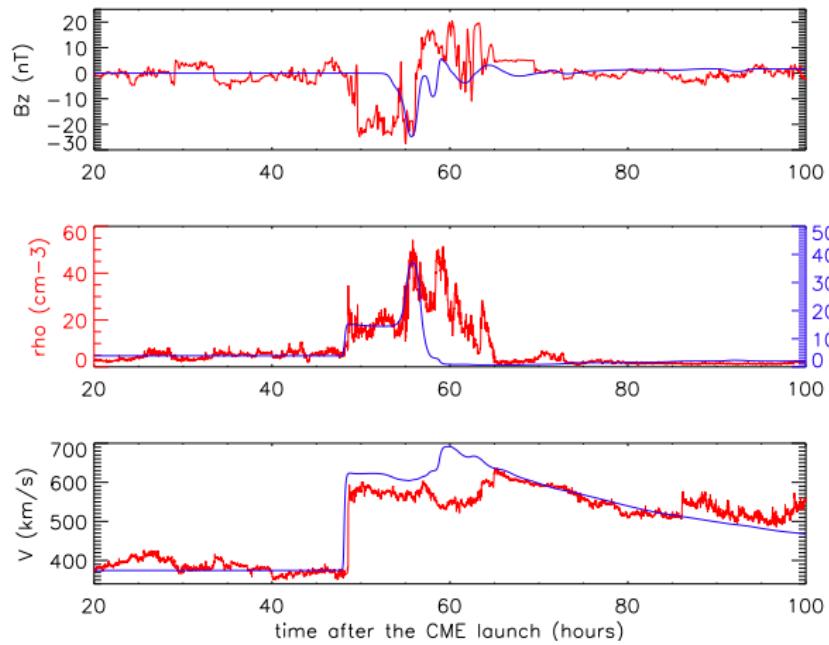
- Normal (blue) and inverse (red) CME
- 3-part structure of CME
 - 1 leading shock
 - 2 dark cavity
 - 3 high density core in cavity
- leading shock front
- . . .



Event study

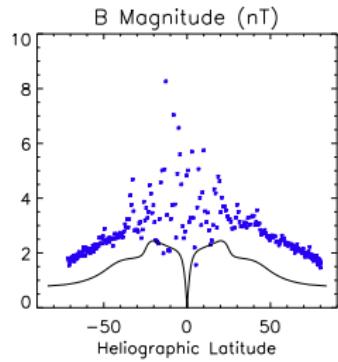
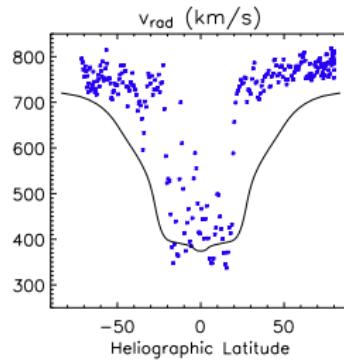
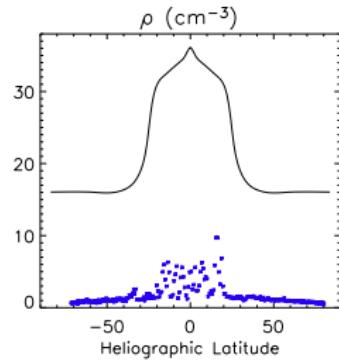
- full halo CME observed by LASCO and EIT on April 4, 2000
- observed at 16:32 UT in C2 frame
- related flare observed by EIT at 15:24 UT
- C3 measurements : plane-of-the-sky speed is 984 km/s
- try to match ACE data by
 - using wind model 2
 - playing with CME parameters (v_{cme} , θ_{cme} , B_{rope} , polarity)

Event study



Best match :
inverse CME
with $v_{\text{cme}} =$
1700 km/s

However...



Comparison of Ulysses data (between 1/11/94 and 1/08/95, i.e. when spacecraft was evolving between 1.34 and 2.03 AU) and wind model 2 at 1 AU

Conclusions

The chosen **background wind model influences** :

- the initiation of the CME (threshold, energetics,...)
 - time of formation (threshold), energetics, speed, acceleration,...
- evolution of the CME
 - shape of leading shock front, shock speed, spread angle, mass distribution,...

Clearly, the **initial parameters** (shear velocity, polarity of fluxrope, v_{CME} , ρ_{CME} , θ_{CME} ,...) also influence the structure and evolution of the CME.

We quantified these effects and are using this knowledge for “**CME SHOCK WAVE seismology**”